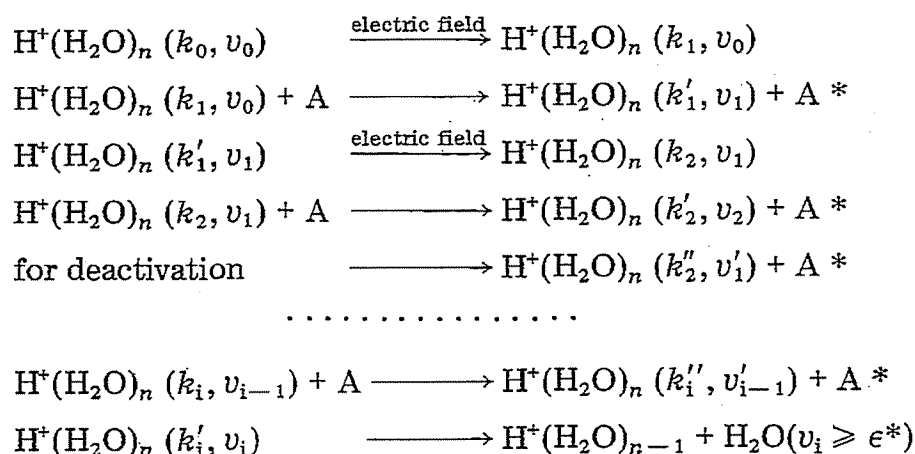


On the other hand, the third mechanism becomes viable when the pressure in the collision region is high. Whether dissociation is caused by single or multiple collision processes depends on the collision energy of the cluster ions and the neutral molecules. The collision energy is determined by the electric field strength and the cluster ion mean free path. Assuming the cluster ions are in their ground energy states, the collision energy must be above the dissociation energy level for cluster ions to dissociate through a single collision.

The mean free path of ions is estimated to be ~ 0.05 mm at ~ 1 torr. Therefore, the average kinetic energy which cluster ions obtain from the electric field in the collision region during a collision interval is estimated to be about 0.1 eV when the electric field is 20 V cm^{-1} . This estimated energy is smaller than the cluster dissociation energy which indicates that dissociation requires multiple collisions.

The probable dissociation process is as follows



where k_i and v_i denote the kinetic and vibrational energies of the cluster ions, respectively, A is any colliding neutral molecule and the asterisk indicates that the particle has excess energy. The kinetic energy of cluster ions after a collision is represented by k'_i . Vibrational and kinetic energies after a deactivational collision are represented by $v'_i (< v_i)$ and k''_i , respectively.

Cluster ions produced at atmospheric pressure enter the collision region with a kinetic energy k_0 and vibrational energy v_0 . The cluster ions are accelerated by the electric field to a kinetic energy k_1 and collide with a neutral molecule A. Part of the kinetic energy k_1 is converted to internal energy in the cluster ion ($v_0 \rightarrow v_1$). Another part is transferred to A. The increase in internal energy depends on the collisional pattern of the cluster ions. The vibrationally-excited cluster ion is again accelerated in the electric field to kinetic energy k_2 . The ion collides again with a neutral molecule, converting its kinetic energy into the internal energy and kinetic energy of the neutral molecule.

When the internal energy becomes large, a deactivation process, in which the internal energy (v_i) is given to the neutral molecule, can not be neglected.

However cluster ions obtain some energy through the collision process as a whole. This is because the internal energy of the cluster ions is spread over all degrees of freedom. Consequently, the kinetic energy which cluster ions obtain from the electric field is repeatedly converted to vibrational energy through collisions. The cluster ion finally dissociates when the resulting internal energy reaches the critical energy level and enough energy is concentrated on a certain vibration corresponding to the dissociation mode.

The average vibrational energy of cluster ions is lower than the dissociation energy D just before dissociation. For example, most of the cluster ions, $H^+(H_2O)_4$, dissociate in a field of 35 V cm^{-1} which corresponds to a potential difference of 14 V between the two aperture electrodes. Therefore the largest internal energy which the ions obtain is 14 eV. If the energy is spread over all 33 internal vibrations, the average vibrational energy increase per degree of freedom is about 0.4 eV. This is smaller than the dissociation energy, 17 kcal mol^{-1} (0.7 eV).

The dissociation of cluster ions occurs when a large amount of energy is concentrated in a certain vibration or when vibrationally-excited cluster ions collide with neutral molecules. After dissociation, the cluster ions are again accelerated in the electric field and collide with molecules increasing their vibrational energies. This sequence continues until the ions enter the analyzing region through the second aperture.

CONCLUSION

The collision-induced dissociation of water cluster ions was investigated at high pressure. The dissociation energy of cluster ions was found to be closely related to the electric field strength, pressure and length of the collision region. The critical electric field strength is also related to the internal energy of ions produced in an API ion source.

It was shown that cluster ions can be reduced to molecular or quasi-molecular ions by means of a high pressure collision-induced dissociation technique. This permits qualitative or quantitative study of an API.

The dissociation mechanism was deduced to depend on the vibrational excitation of cluster ions through many collisions. The collision-induced dissociation process at high pressure is related to the ion-molecule reaction, energy relaxation, cluster ion structure, collisional energy transfer, unimolecular reaction, etc. Thus this collisional dissociation study is applicable to studies of the molecular structure of bond dissociation energy of cluster ions. It can also be applied to studies of molecular structure or bond dissociation energy of organic molecular ions by increasing the electric field strength to dissociate the molecular ions. However, further precision regarding the mechanism will entail many experiments.

ACKNOWLEDGEMENT

The authors express their thanks to Dr. M. Yamane and Dr. S. Yamamoto for helpful discussions.

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US005298743A

United States Patent [19][11] **Patent Number:** **5,298,743****Kato**[45] **Date of Patent:** **Mar. 29, 1994**[54] **MASS SPECTROMETRY AND MASS SPECTROMETER**[75] **Inventor:** Yoshiaki Kato, Mito, Japan[73] **Assignee:** Hitachi, Ltd., Tokyo, Japan[21] **Appl. No.:** 942,992[22] **Filed:** Sep. 10, 1992[30] **Foreign Application Priority Data**

Sep. 12, 1991 [JP] Japan 3-232956

[51] **Int. Cl.⁵** B01D 59/44; H01J 49/00[52] **U.S. Cl.** 250/288; 250/289[58] **Field of Search** 250/288, 288 A, 289[56] **References Cited****U.S. PATENT DOCUMENTS**

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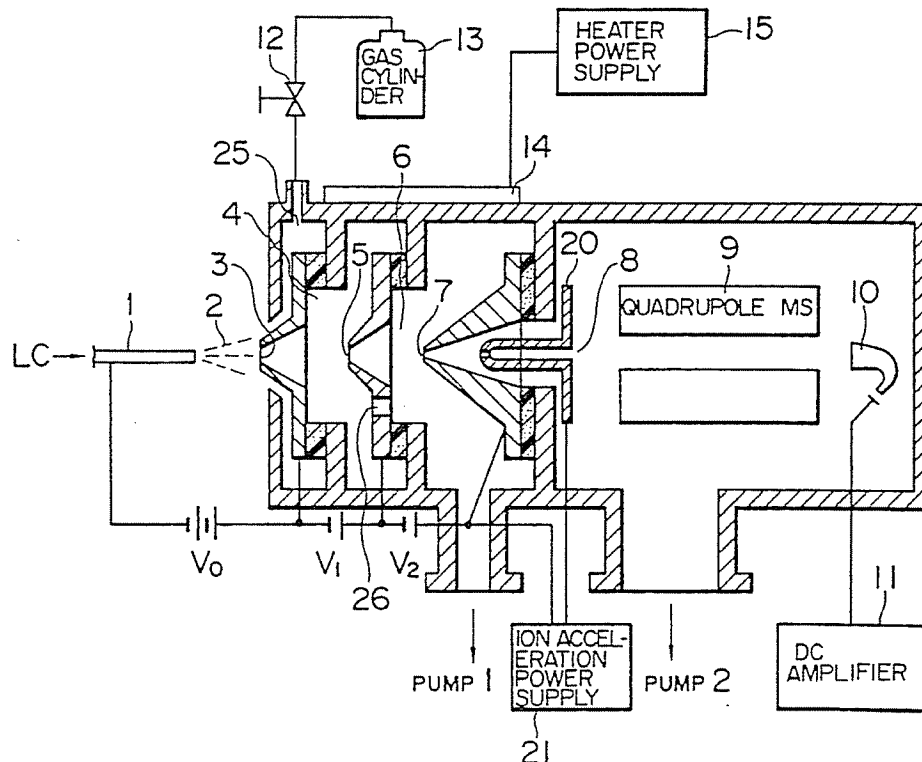
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et al: "Atmospheric pressure ionization mass spectrometry", p. 713A; figures 1-2.

Primary Examiner—Bruce C. Anderson
Attorney, Agent, or Firm—Antonelli, Terry, Stout & Kraus[57] **ABSTRACT**

Ions generated under an atmospheric pressure pass through vacuum chambers partitioned through first, second and third fine holes. The ions are led to an MS part where the ions are mass-analyzed. A first vacuum chamber adjacent to an atmospheric pressure part has not vacuum pump for independently pumping this chamber. The first vacuum chamber is evacuated by a common pump together with a second vacuum chamber via a bypass hole formed in the wall having the second aperture. A pressure of the first vacuum chamber can be set to several 100 Pa, while a pressure of the second vacuum chamber can be set to several 10 Pa. Sufficient desolvation has been attained by an ion acceleration voltage of approximately 100 V in the first vacuum chamber, while a speed spread can be restrained. The ions are accelerated by approximately 10 V in the second vacuum chamber, and the speed spread can be restrained as low as possible.

16 Claims, 8 Drawing Sheets

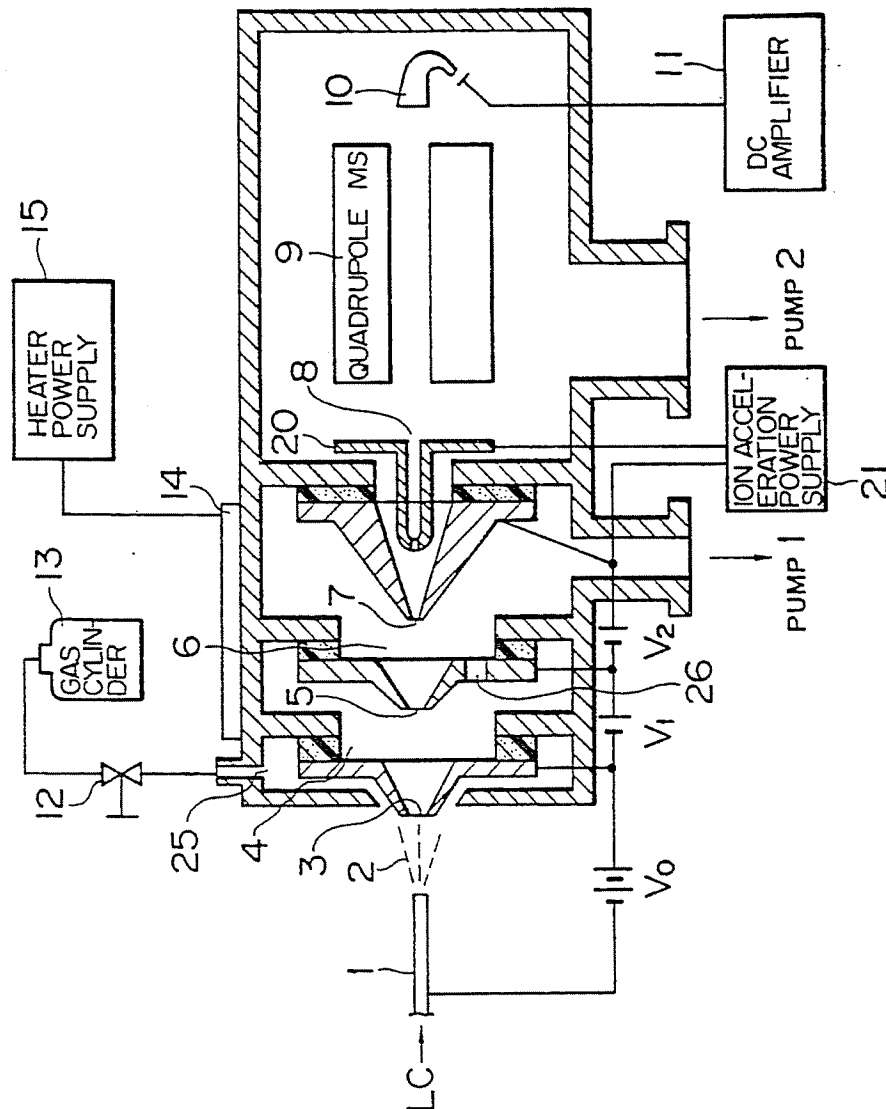
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FIG. 1



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FIG. 2

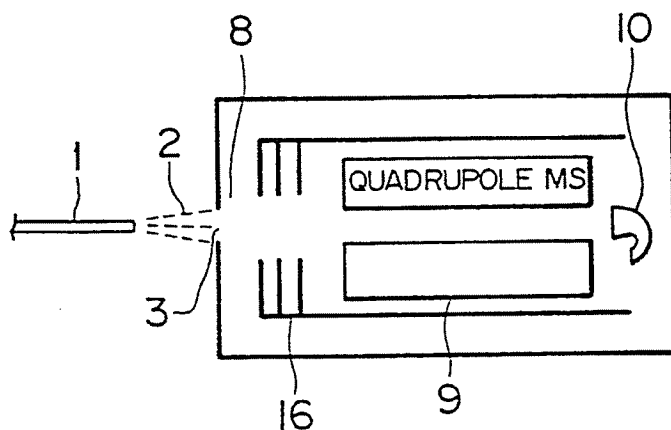


FIG. 3

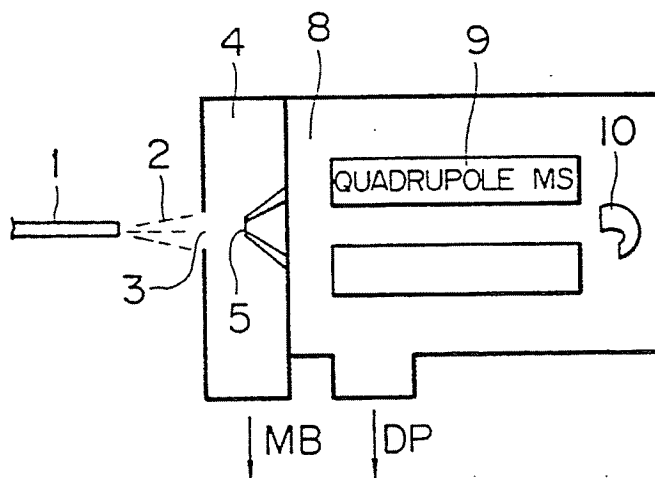


FIG. 4

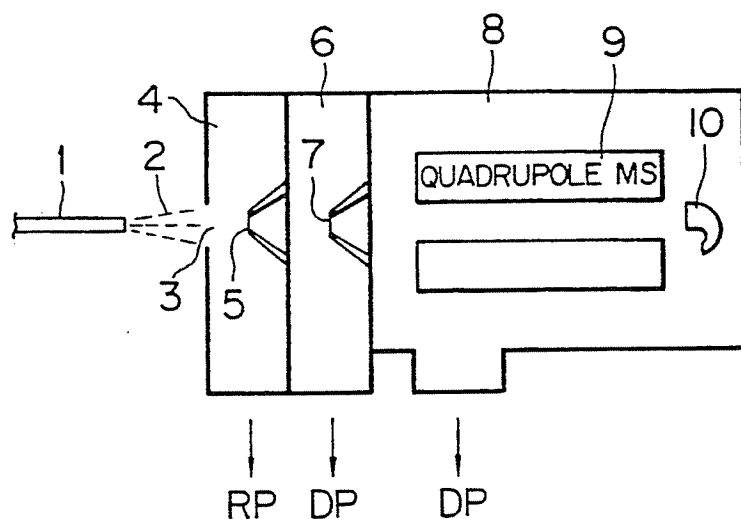


FIG. 5

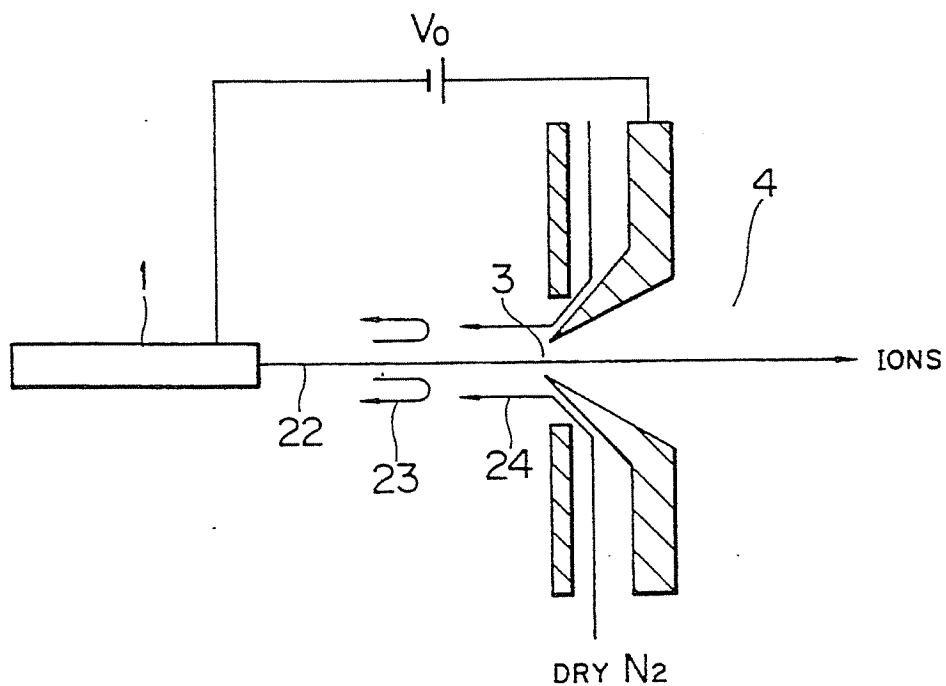


FIG. 6

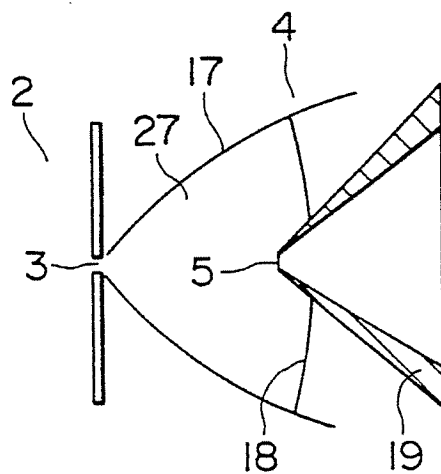


FIG. 7

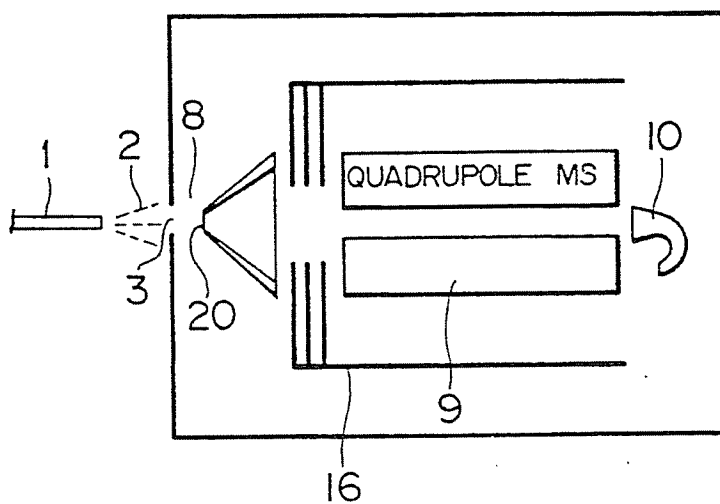


FIG. 8

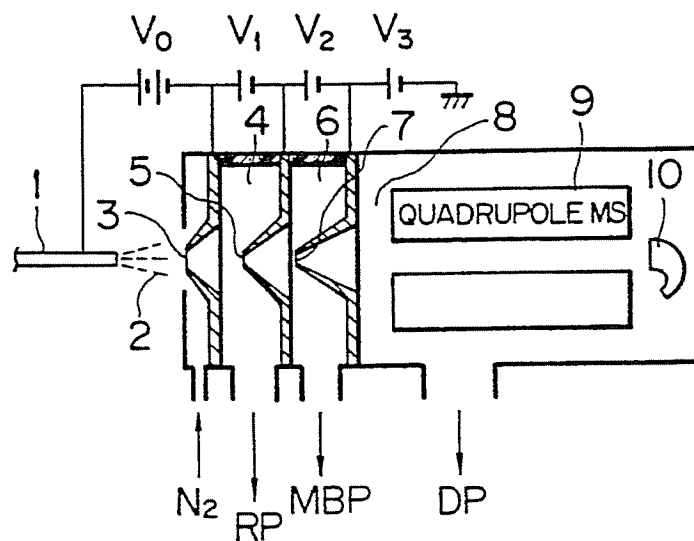


FIG. 9

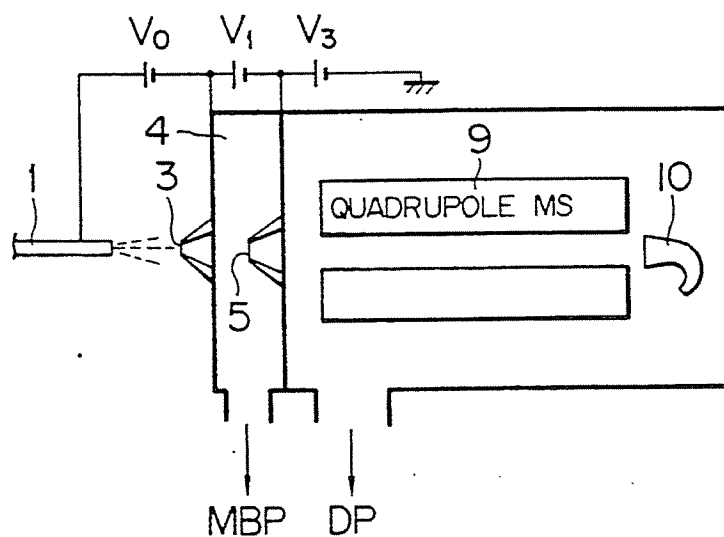


FIG. 10

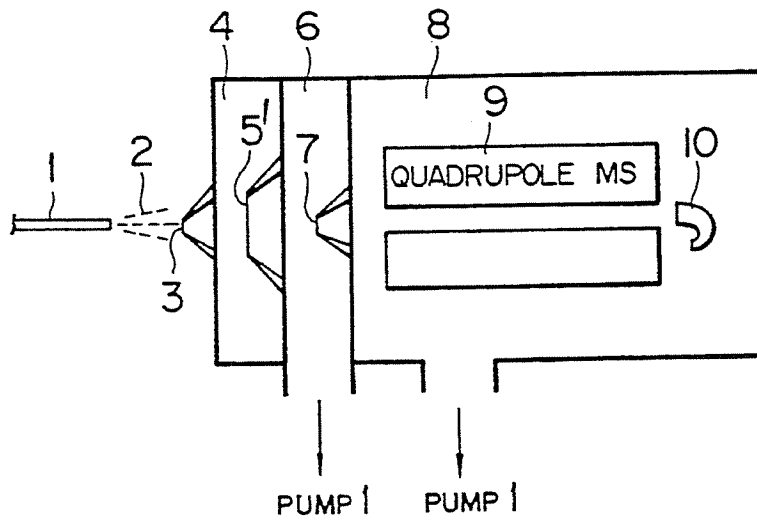
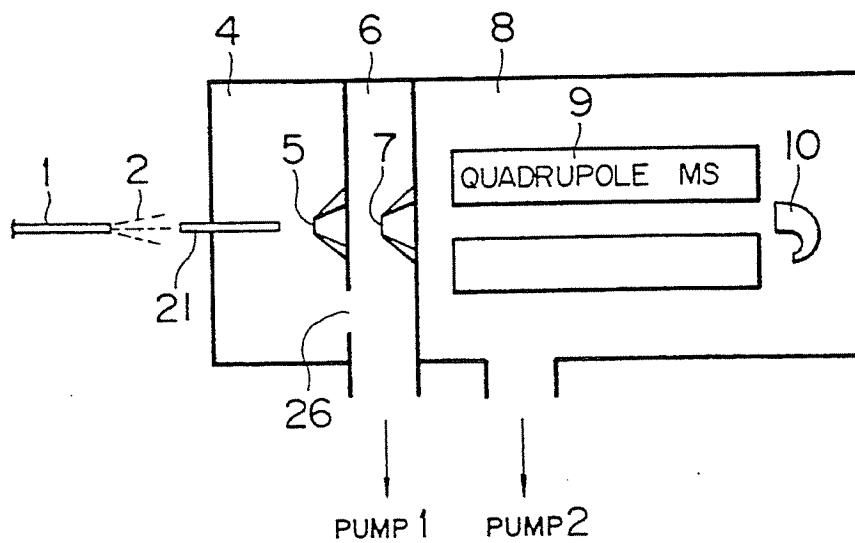


FIG. 11



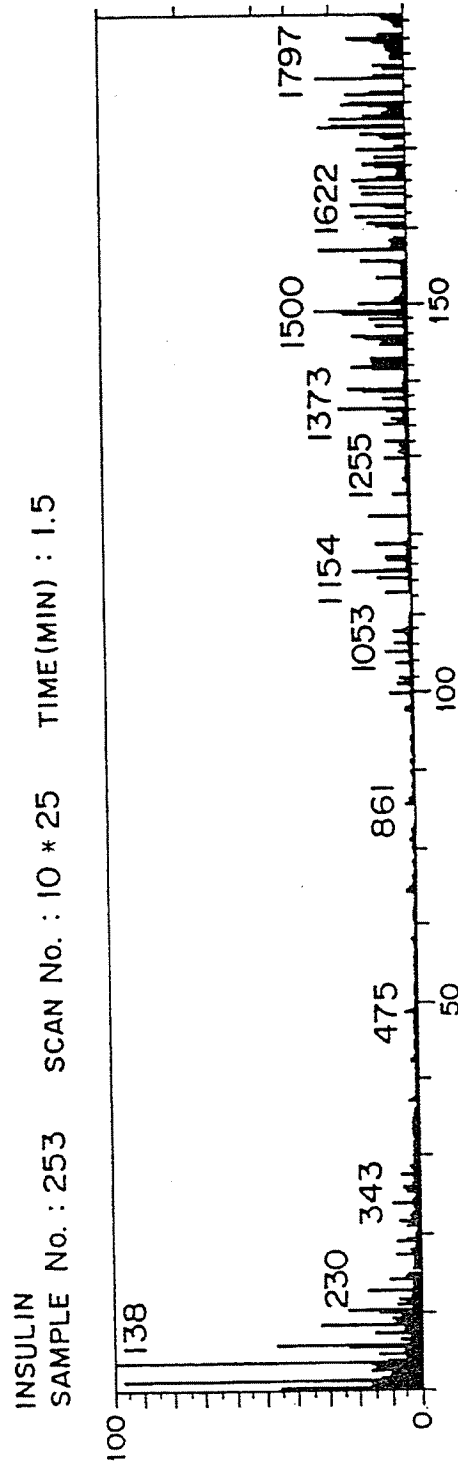
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FIG. 12



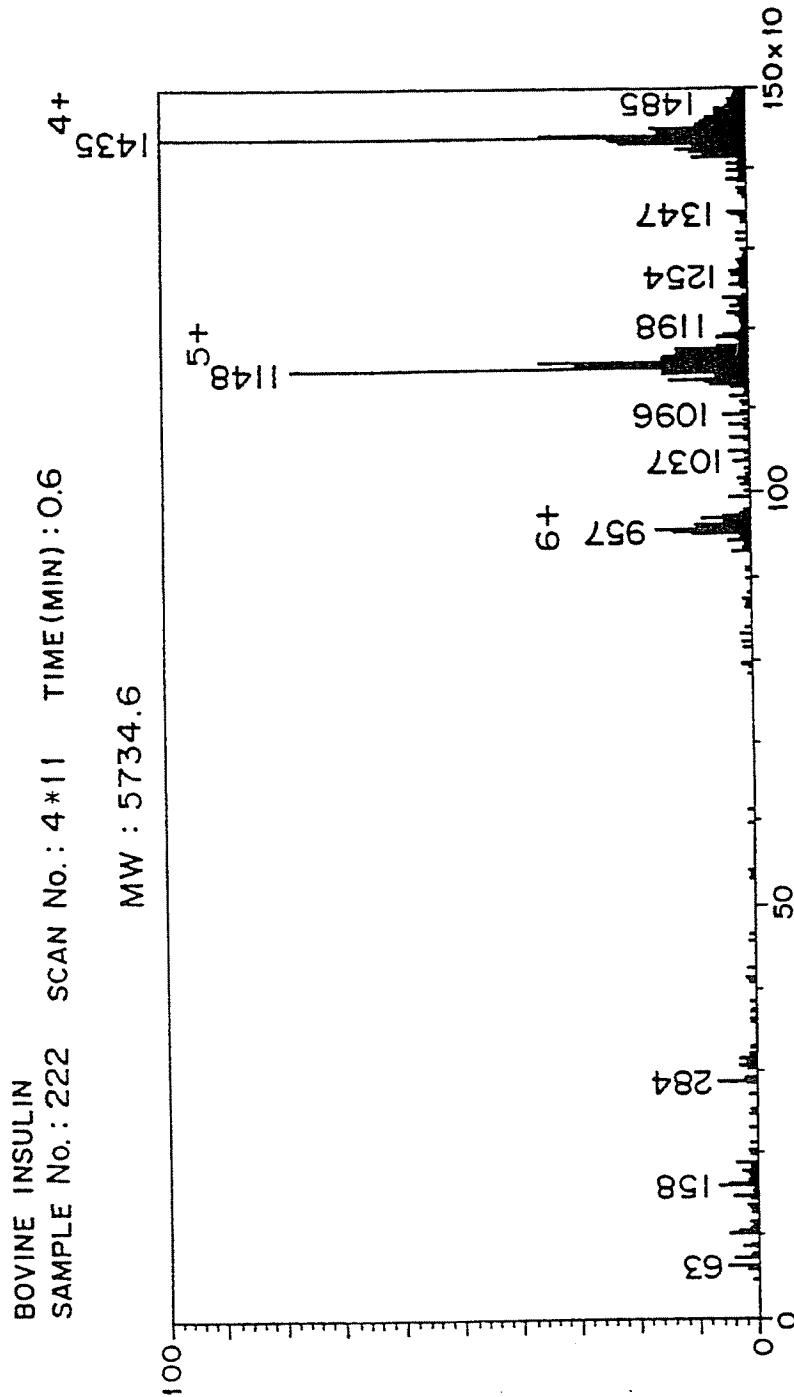
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FIG. 13



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MASS SPECTROMETRY AND MASS SPECTROMETER

FIELD OF THE INVENTION

The present invention relates generally to a mass spectrometry (method of mass analysis) and mass spectrometer (apparatus for mass analysis) and, more particularly, to a mass spectrometry and mass spectrometer for generating ions under an atmospheric pressure and analyzing masses.

DESCRIPTION OF THE RELATED ARTS

Atmospheric pressure ionization (API) is often utilized for mass-analyzing a fluid containing sample and solvent components flowing from a liquid chromatograph (LC). In this atmospheric pressure ionization, soft ionization is effected so as not to impart an excessive energy to sample molecules. For this reason, the sample is decomposed to a less extent upon ionization, and the molecular ions are easy to observe. Further, because of the ionization under a high pressure (atmospheric pressure), even a substance having a low ionization potential is ionized at a high ionization efficiency. Therefore, a highly sensitive mass analysis can be expected. The ionization under an atmospheric pressure is described in detail in Analytical Chemistry, Vol. 62, No. 13, pp. 713A-725A (1990).

The ions have to be introduced into a vacuum in order to mass-analyze the ions generated under the atmospheric pressure. If the ions generated under the atmospheric pressure are immediately led into a high vacuum chamber to perform the mass analysis, there arises problems such as contamination in the high vacuum chamber. Hence, in most of the cases, low and intermediate vacuum chambers are provided between the atmospheric pressure and the high vacuum to give a gradual pressure gradient between the atmospheric pressure and the high vacuum, while these chambers are evacuated independently by use of vacuum exhaust pumps.

However, in the case of differentially evacuating the low and intermediate vacuum chambers in that way by use of the independent separate vacuum systems, the vacuum systems become complicated and expensive.

SUMMARY OF THE INVENTION

It is an object of the present invention to provide a mass spectrometry and mass spectrometer capable of simplifying the vacuum systems.

According to the present invention, low and intermediate vacuum chambers are provided between an atmospheric pressure ionizing unit and a high vacuum unit for effecting a mass analysis and are evacuated by a common vacuum system.

According to the present invention, the low and intermediate vacuum-chambers are evacuated in this way by the common vacuum system, and hence the vacuum system is simplified. This in turn leads to a reduction in costs.

The foregoing and other objects, features as well as advantages of the invention will be made clearer from the description of preferred embodiments hereafter referring to the attached drawings.

BRIEF DESCRIPTION OF THE DRAWINGS

FIG. 1 is a schematic diagram of a whole arrangement of a liquid chromatograph/mass spectrometer,

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showing one embodiment according to the present invention;

FIG. 2 is a conceptual diagram illustrating an LC/MS device based on the conventional technique;

FIG. 3 is a conceptual diagram illustrating the LC/MS device based on the conventional technique;

FIG. 4 is a conceptual diagram illustrating the LC/MS device based on the conventional technique;

FIG. 5 is a conceptual diagram depicting the LC/MS device based on ionization in a counter gas system;

FIG. 6 is a schematic diagram showing a shock wave by a supersonic fluid introduced into a vacuum from an atmospheric pressure;

FIG. 7 is a conceptual diagram illustrating the LC/MS device including an ion acceleration electrode for restraining a spread of speed.

FIG. 8 is a conceptual diagram of the LC/MS of a 3-stage differential pumping system;

FIG. 9 is a conceptual diagram of the LC/MS of a 2-stage differential pumping system;

FIG. 10 is a conceptual diagram of the LC/MS device, showing another embodiment of the present invention;

FIG. 11 is a conceptual diagram of the LC/MS device, showing still another embodiment of the present invention;

FIG. 12 is a diagram showing an insulin mass spectrum obtained by the conventional system; and

FIG. 13 is a diagram showing the insulin mass spectrum obtained in accordance with the embodiment of the present invention.

DETAILED DESCRIPTION OF THE PREFERRED EMBODIMENTS

In advance of describing embodiments of the present invention, the background and fundamentals of the present invention will be at first explained.

For mass-analyzing ions generated under an atmospheric pressure, at first, the ions have to be introduced into a vacuum. Further, for a high sensitivity measurement, it is required that the ions be led to a high vacuum mass spectrometer (MS) so as to minimize a loss of the ions generated under the atmospheric pressure (at a high efficiency). For this purpose, vacuum system is the first item which has to be considered in an LC/MS interface, i.e., a mass spectrometer directly connected to the liquid chromatograph (LC). Thus vacuum system is classified roughly into two systems. The first system is, as illustrated in FIG. 2, a method of partitioning an atmospheric pressure part 2 and a vacuum part 8 by use of a partition wall formed with an aperture 3 and sampling the ions generated via this aperture 3. The second system is, as depicted in FIG. 3 or 4, a method of introducing the ions to an MS part 8 through several-staged differential pumping systems employing a plurality of partition walls formed with the aperture 3 and skimmer(s) 5, 7. In the first system, an aperture diameter d (m) and a pumping speed S (m^3/s) of a vacuum pump are given as follows to obtain a vacuum required for the MS. A vacuum degree for operation of the MS is herein 10^{-3} - 10^{-4} Pa. A conductance C_1 of a gas in viscous flow region of the aperture diameter d (m) is obtained by the formula (1).

$$C_1 \approx 158d^2$$

(1)

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Assuming that the pumping speed of the vacuum pump for the MS part is S_1 (m^3/s), the vacuum P_1 of the MS part is obtained by the formula (2). Besides, the atmospheric pressure P_0 is approximately 10^5 Pa.

$$P_1 = C_1(P_0 - P_1)/S_1 \quad (2)$$

The formula (2) can, because of $P_0 \gg P_1$, be approximate to the formula (2').

$$P_1 = C_1 \cdot P_0 / S_1 \quad (2')$$

$$= C_1 \times 10^5 / S_1 \quad (3)$$

Now, assuming that the vacuum pump for the MS part is an oil diffusion pump having a pumping speed of 1,000 liters/s = $1 \text{ m}^3/\text{s}$, the aperture diameter d required for accomplishing a vacuum degree of 10^{-4} Pa in the MS part is given as follows. From the formulae (1) and (3),

$$\begin{aligned} d &= \sqrt{\{P_1 \cdot S_1 / (157 \times 10^5)\}} \\ &= \sqrt{\{10^{-4} \times 1 / (157 \times 10^5)\}} \\ &= 2.52 \times 10^{-6} \text{ (m)} \end{aligned} \quad (4)$$

Namely, the aperture has a diameter of approximately $2.5 \mu\text{m}$. If a cryopump having a pumping speed of 10,000 liters/s is employed as the vacuum pump, the aperture diameter is nothing more than $7.9 \mu\text{m}$. When the ions are sampled from the atmosphere through the aperture having such a small diameter clogging of the aperture is frequently caused due to matters such as dusts in the air. Further, since the diameter of the aperture is small, a good deal of ions can not be introduced. This makes the high sensitivity measurement difficult. An additional problem is that the cryopump is remarkably expensive. FIG. 2 is a schematic diagram based on this system. The ions sprayed from a spray nozzle 1 and generated under an atmospheric pressure and in a high electrostatic field 2 enter the MS part via the aperture 3. Neutral molecules are trapped by a cooling fin 16 of the cryopump. On the other hand, the ions go straight and undergo a mass sorting in a quadrupole MS 9 and reach a detector 10.

In the case of a system (FIG. 3) based not on such an arrangement that the ions are sampled directly through the single aperture but on such an arrangement that two or more apertures are disposed in series on the same axis; and vacuum regions between partition walls each having therein the aperture is performed by independent vacuum pumps, the vacuum of the MS unit 8 is defined as follows. Let P_0 be the atmospheric pressure, and let P_2 be the vacuum degree of the MS unit 8. Let S_1 , S_2 be the pumping speeds of the vacuum pumps of the differential pumping system part and MS part, respectively. Let C_1 , C_2 be the conductances of gases of the first and second apertures 3, 5, respectively. Further, let d_1 , d_2 be the diameters of the first and second apertures.

The pressure P_1 of the differential vacuum chamber 4 is given by the following formula:

$$\begin{aligned} P_1 &= C_1(P_0 - P_1)/S_1 \\ &\approx C_1 P_0 / S_1 \end{aligned} \quad (5)$$

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Besides, the pressure P_2 of the MS part is given by the following formula:

$$\begin{aligned} P_2 &= C_2(P_1 - P_2)/S_2 \\ &\approx C_2 \cdot P_1 / S_2 \end{aligned} \quad (6)$$

It is because $P_0 \gg P_1 \gg P_2$. From the formulae (5) and (6),

$$P_2 = C_1 \cdot C_2 P_0 / (S_1 S_2) \quad (7)$$

is derived.

Further, C_1 is given by the formula (1).

$$C_1 = 157 d_1^2 \quad (8)$$

The conductance C_2 in the molecular flow region is given by the following formula:

$$C_2 = 116 \times A \quad (9)$$

However, A is the area of the aperture. This is further expressed as:

$$C_2 = 116 \times \pi (d_2/2)^2 \quad (10)$$

Hence, the formula (7) is expressed as:

$$P_2 = 157 d_1^2 \times 116 \times \pi (d_2/2)^2 P_0 / (S_1 \times S_2) \quad (11)$$

Now, it is assumed that the MS part 8 is evacuated by the oil diffusion pump having a pumping speed of 1,000 liters/s, while the differential vacuum system part 4 is evacuated by a mechanical booster pump of 16.7 liters/s. The diameter of the first aperture 3 is assumed to be $200 \mu\text{m}$, while the diameter of the second aperture 5 is assumed to be $400 \mu\text{m}$. The vacuums P_1 , P_2 of the differential pumping system part 4 and MS part 8 are respectively given from the formulae (5) and (11):

$$P_1 = 37.6 \text{ (Pa)}$$

$$P_2 = 5.6 \times 10^{-4} \text{ (Pa)} \quad (12)$$

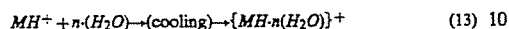
The vacuum of the MS part 8 is high enough for the mass analysis. As compared with the first system employing the single aperture and the high speed vacuum pump, the second system using a plurality of apertures and the differential pumping system exhibits such an advantage that the large apertures and the inexpensive vacuum pumps can be utilized. For this reason, the second system is widely utilized in a great number of vacuum devices. Further, as illustrated in FIG. 4, a 3-stage differential pumping system is similarly utilized. This differential pumping system corresponds to a method which is excellent in terms of such a point that the ions generated under the atmospheric pressure are led to the MS part at a high efficiency. In general, the two-or three-stage differential pumping system is used in the LC/MS.

There is also a point to be considered other than the vacuum in the LC/MS interface. When the ions generated under the atmospheric pressure are introduced into the vacuum, a rapid adiabatic expansion takes place. Thus, the introduced ions and molecules are rapidly cooled off. Therefore, the molecules such as those of water and alcohol which have been introduced together with the ions into the vacuum are added to the ions,

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resulting in a generation of cluster ions. Especially in the case where the sample ion has a good number of charges, or where the ion has a multiplicity of functional groups with a high polarity, there are generated the cluster ions in each of which many molecules such as molecules of water and alcohol are added to the ion. For instance, in the case of an addition of water, this is expressed by the following formula:



The cluster ion is an ion to which a multiplicity of polar molecules are added. However, the type and the number of the molecules to be added are not constant. It is therefore impossible to directly obtain the information on a molecular weight of the sample molecule from the cluster ion by means of the MS. Further, ions having same m/z are distributed widely in the form of a multiplicity of cluster ions, and hence a detected ionic current value is also decreased. Therefore, desolvation for removing the added molecules from the cluster ions is required. Proposed as a method therefor are the following methods and a combinational system thereof. In any case, an external energy greater than the addition energy of the polar molecules is given to the cluster ion, thereby releasing the polar molecules from the ion. If the externally given energy is excessive, the cluster ions are decomposed, and molecular weight information can not be given. Whereas if too low, the release of the added molecules is insufficient, and molecular weight information can not be given either. Therefore, the energy imparted to the cluster ion is controlled to exceed slightly the energy that is required for the release of the added molecules. It is required that the energy be repeatedly injected into the ions.

For release of neutral polar molecules, there may be several possible measures as follows;

- (1) Collision with counter gas
- (2) Adiabatic compression on Mach disk surface
- (3) Heating
- (4) Ion acceleration and collision

(1) Collision with Counter Gas

FIG. 5 is a schematic diagram showing this system. The cluster ions are made to pass through an inert gas which has been heated ($\sim 70^\circ \text{C}$), e.g., dry nitrogen. Nitrogen molecules are caused to collide with the cluster ions, and the heat is transferred to the cluster ions from the nitrogen molecules continually, thereby releasing the added molecules from the ions. The dry nitrogen is flowed in a direction 24 opposite to a flow of the ions in the vicinity of the ion sampling aperture 3. Therefore, neutral solvent molecules (such as water) flowing together with the ions are flowed back in a direction 23 opposite to the ions sampling aperture 3 due to the dry nitrogen. On the other hand, the ions 22 are accelerated by an electric potential applied between the aperture 3 and the spray nozzle 1 and collide with the dry nitrogen molecules. The ions 22 undergo the desolvation and enter the aperture 3. This also prevents extra polar molecule from entering the vacuum chamber, and a possibility of collision and recoupling within the vacuum chamber can be made low. Although the perfect desolvation is not attainable only by the collision with the counter gas, this system is a preferable method capable of restricting the polar molecules from entering the vacuum chamber. Hence, the desolvation is

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attainable more efficiently in a combination with the following system than used singly.

(2) Adiabatic Compression of Mach Disk Surface

Gaseous molecules having entered via the aperture from the atmospheric pressure are changed into a supersonic flow of molecules. Consequently, as illustrated in FIG. 6, a Mach disk 18 and a barrel shock 17 depending on the pressure in the vacuum chamber are produced. Where P_0 is the pressure of the outside of the aperture 3; P_1 is the pressure in the vacuum chamber 4; and d is the aperture diameter, the Mach disk is generated on a distance X_M from the aperture 4.

$$X_M = (2/3)d \sqrt{(P_0/P_1)} \quad (14)$$

For example, assuming that the pressures in front and in rear of the aperture having a diameter of 0.3 mm are 10^5 Pa (atmospheric pressure) and 100 Pa, the Mach disk is expressed as:

$$\begin{aligned} X_M &= (2/3) \times 0.3 \times \sqrt{(10^5/10^2)} \\ &= 0.2 \times 31.6 = 6.3 \text{ (mm)} \end{aligned} \quad (15)$$

Namely, the Mach disk is generated in a place positioned 6.3 mm away from the aperture towards the high vacuum part. The adiabatic compression is effected on the Mach disk surface, whereby the cluster ions are rapidly heated. As a result, the desolvation is performed. Where the second aperture 5 is disposed in a place positioned 7 mm or more apart backwards from the first aperture 3, the cluster ions invariably pass through the Mach disk surface, thereby promoting the desolvation with heating by the adiabatic compression. This system is a preferable method capable of attaining the desolvation without supply of special external energy. In rear of the Mach disk, however, the flow of molecules becomes absolutely irregular, and the flow of ions entering the second aperture does not become constant. This causes such a defect that a sampling yield of the ions does not increase. Generally, for improving the ions sampling yield, sampling is often effected in a molecular flow region (Silent Zone) 27 in front of the Mach disk where the ions and gas molecules continue their motion in straight line. However, if sampling is effected in the molecular flow region 27, as a matter of course, the desolvation and the adiabatic compression by the Mach disk are not carried out. This implies that a well-directed flow of abundant molecules are merely sampled.

(3) Heating

The gas diffused into the vacuum from the atmospheric pressure is rapidly cooled by the adiabatic expansion. In a case where the gas to be introduced is heated beforehand, and where the interface including the aperture is heated, the adiabatic cooling can be compensated to some extent, and an addition of water and the like can be prevented. It is, however, difficult to attain the perfect desolvation only by heating. It is because most of ions of organic compounds passing through this interface tend to easily undergo the thermal decomposition by heating. It is therefore impossible to perform heating at a high temperature for the purpose of the desolvation.

(4) Ion Acceleration and Collision

If the pressure reaches 100 Pa–10 Pa, a mean free path of the gaseous molecules become about 0.06 to 0.6 mm. When an electric field is applied under such a pressure, the ions existing in the gas are accelerated in a direction along the electric field and collide with the neutral molecules. During a flight of the ions in the electric field, the acceleration and collision are repeated. When the mean free path is 0.1 mm (~66 Pa), the ions are accelerated by approximately 1 eV in the electric field of 100 V/cm, where e is the number of electric valences of the ions. A part of this kinetic energy is transformed into an internal energy (thermal energy) by the collision. If a value of this internal energy exceeds the addition energy (several kJ/mol–several 10 kJ/mol=0.01 eV–0.1 eV) of the molecules of water and the like, the water molecules etc. can be released. Important factors in this desolvation system are a vacuum degree and an intensity of the electric field in the case of the acceleration and collision. Generally, as illustrated in FIG. 8, the electric potential is applied between the first and second apertures 3, 5 or/and between the second and third apertures 5, 7, whereby the ions are accelerated and collide with the neutral molecules. A degree of the desolvation can be changed by controlling the applied voltages V_1 , V_2 . This method is remarkably effective in the desolvation. This method, however, has a defect of directly undergoing influences of the pressures of the ion acceleration and collision parts 4, 6. Besides, because of accelerating the ions, there is a risk in which a part of the kinetic energy is not consumed by the collision but is imparted directly to the ions. Therefore, the ions which have entered the high vacuum MS part 8 spread in speed. It follows that this directly brings about declines in resolving power and sensitivity in the mass analysis. If the speed spread exceeds 1 eV, it is difficult to attain the resolving power more than one mass unit in the case of the quadrupole MS. In addition, a transmissivity of the ions is also decreased. In the case of a double focusing mass spectrometer, the large energy dispersion occurs due to the electric field, with the result that the declines in the sensitivity and resolving power are induced.

The mean free path of the nitrogen molecules under from the atmospheric pressure (~ 10^5 Pa) to 10^3 Pa is approximately 5×10^{-5} mm– 5×10^{-3} mm. Even when the electric field of 100 V/mm is applied under these pressures, the kinetic energy received by the ions ranges from 5×10^{-3} eV to 5×10^{-1} eV, which is considerably lower than 1 eV. The collisions frequently happen in this pressure region, and it is therefore impossible to accelerate the ions, although the ion moving direction can be changed even when the electric field is applied. More specifically, even when the ions are accelerated under this pressure, the spread of the kinetic energy can be restrained not more than 1 eV. On the other hand, under 10^3 Pa through 1 Pa, the mean free path of the nitrogen molecules ranges from approximately 5×10^{-3} mm to 5 mm. When the electric field of 100 V/mm is applied under this pressure, the kinetic energy received by the ions within the mean free path is as large as 5×10^{-1} eV to 5×10^2 eV. This causes a large spread of the kinetic energy (speed). On the other hand, in the vacuum of 0.1 Pa to 10^{-4} Pa, the mean free path becomes 50 mm to 50 m. Reduced is a probability that the accelerated ions collide with the neutral molecules in the acceleration field. The spread of the kinetic energy

is reduced. On the occasion of effecting the ion acceleration and a dissociation of collision, it is necessary to consider this spread of the kinetic energy together. As described above, if the ions are accelerated in the low vacuum (10^3 Pa or more) or in the high vacuum (10^{-1} Pa or less), the spread of the speeds of the ions is negligible. There have been already described the advantages in terms of the vacuum system based on the system which utilizes the differential pumping system to take the ions, generated under the atmospheric pressure, into the high vacuum MS. The ions are converged by applying the electric potential between the apertures of this differential pumping system and can be highly efficiently introduced into the MS. Further, at the same time the desolvation by the acceleration, collision and dissociation can be effected. However, the creation of spread of the ion speeds in the process of this desolvation gives an adverse effect.

In the case of the system, shown in FIG. 7, for taking the ions directly from the atmospheric pressure into the MS part, the vacuum gradually becomes higher from the ions sampling aperture 3 in the ion flying direction of the MS part 8. If there is a sufficient space between the ions sampling aperture 3 and an ion acceleration electrode 20, the ions are accelerated between these two portions and invariably pass through the intermediate pressure region (10^3 Pa–1 Pa). Spread of energies of the ions do not occur in the high pressure part (10^5 – 10^3 Pa). On the other hand, the ions are accelerated in the region where the pressure ranges from 10^2 Pa to 1 Pa, and the energy spread is provided. In order to restrain the energy (speed) spread as low as possible, the ion acceleration electrode 20 is positioned close to the ion sampling aperture 3, and the ions are accelerated in the high pressure part (10^5 – 10^3 Pa). In this region, however, the cluster ions can not be sufficiently accelerated. The energy required for the desolvation cannot be given to the cluster ions. Therefore, the desolvation in this region can not be expected.

In the case of the differential pumping system of FIG. 3 also, the ion acceleration in the differential pumping system part is an acceleration in the intermediate pressure region (10^3 –1 Pa), and it follows that the energy spread is imparted. The following prevention measures are required for avoiding this energy spread. The pressure difference is controlled stepwise and accurately by using a plurality of differential pumping system. Further, the desolvation by acceleration is performed in the vacuum of 10^2 Pa or under, and the ion acceleration is restrained at the possible lowest level under the intermediate pressure of 10^2 –1 Pa. The ions are accelerated at a stretch in the next high vacuum region. This requires a difficult of the pressure control and an intricate and expensive differential pumping system as shown in FIG. 8.

Referring to FIG. 8, the pressure of the first vacuum chamber 4 is kept at 10^3 – 10^2 Pa, while the ion acceleration voltage V_1 is kept at 100–200 V. The second vacuum chamber is maintained at 10–1 Pa, while the ion acceleration voltage V_2 is restrained down at 10–20 V. As described above, the collision dissociation is promoted by increasing the ion acceleration voltage in such a low vacuum region as to exert no influence on the ion speed. Whereas in such a region as to exert an influence on the ion speed, the ion acceleration voltage is restrained low. It is not, however, easy to constantly control the pressure and the ion acceleration voltage. Besides, when the high voltage is applied under the

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intermediate pressure (10^3 –1 Pa) for promoting the desolvation, a glow discharge readily starts. Once the glow discharge starts, the ions introduced to the interface disappear. Therefore, the pressure under which the discharge can be avoided and the desolvation can be attained is limited. Typically, 5×10^3 Pa–50 Pa is a pressure suitable for the desolvation.

The present invention is embodied by the following technique.

In the high pressure region (atmospheric pressure 10^5 Pa– 10^3 Pa), the motions of the ions are remarkably restricted even in the electric field. Hence, the control of the direction of the motions of the ions is accomplished by the electric field, and the spread of the speed of the ions is not caused. In the region of 10^2 Pa–1 Pa, the ions are accelerated and repeatedly collide with the neutral molecules. As a result, a large spread of speed of the ions is caused. Further, in the high vacuum of 1 Pa or lower, a probability of collision of the accelerated ions with residual molecules becomes low, and resultantly the speed spread is also decreased. Namely, if the ions are accelerated in the intermediate region (10^2 –1 Pa) between the case of the high pressure and the case of high vacuum, the spread of speed is induced. For this reason, the intermediate vacuum region is physically separated from each of the high-pressure part and high-vacuum part through partition walls with orifices. The voltage required for accelerating the ions is applied in each vacuum chamber. Chambers are provided so that the interface parts are depressurized sequentially from the atmospheric pressure. The chamber adjacent to the atmospheric pressure is evacuated not by an independent pump but through a bypass hole opened to the high vacuum part at the next stage, so that a pressure of this chamber can be easily set by a conductance of this hole. The vacuum pump, the pumping duct and the control power supply of vacuum system can be thereby simplified.

It is easy to keep different chambers under different pressures respectively by a single or common pumping system. The ions are accelerated by the electric field of 200 through 100 V/5 mm in the chamber held at a pressure of 10^3 to 10^2 Pa. It is therefore possible to provide the number of collisions and energy required for the desolvation while restraining the energy spread within 1 eV. An electric potential (approximately 10–20 V/5 mm) enough to converge the ions is given in the chamber of 10^2 to 1 Pa. The energy spread in this region can be thereby restrained within 12 eV.

For describing the embodiment of the present invention with reference to FIG. 1, an ESI (Electro-Spray Ionization, i.e., ionization by spraying a liquid in a high electric field) interface is composed of a spray nozzle 1 to which a high voltage V_0 is applied, a counter gas introduction chamber 25, a first aperture (ion sampling aperture) 3, a first vacuum chamber 4, a second aperture 5, a second vacuum chamber 6, a third aperture 7, an ion acceleration power supply 21, a heater 14 and a heater power supply 15.

An eluate fed in from the LC reaches the spray nozzle 1 and is sprayed in the atmosphere 2. A good deal of electric charges are carried on the sprayed droplet surfaces. The droplets are diminished by evaporating the solvent from the droplet surfaces while flying in the atmosphere 2. When a repulsion of the electric charges of the same polarity carried on the surface becomes greater than a surface tension, the droplets are segmented at a stretch. Finally, it comes to a result that the

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ions have evaporated from the liquid phase to the atmosphere 2 (gas phase). With the intention of helping the segmentation of the droplets and preventing the neutral polar molecules (water, etc.) from entering the interface, the counter gas is made to flow into the atmosphere 2 from the vicinity of the first aperture 3 in a direction opposite to the flying direction of the ions, where the counter gas is fed via a needle valve 12 from a gas cylinder 13. The counter gas is typically heated at 60° – 70° C., thus promoting the evaporation of the solvent from the droplets. The ions move with the aid of the electric field while resisting a flow of the counter gas and enters the first vacuum chamber 4 via the first aperture 3. The ions are then accelerated by the voltage V_1 applied between the partition walls, of the first vacuum chamber, formed respectively with the first and second apertures 3, 5. The ions then collide with the neutral gaseous molecules and undergo the desolvation. The ions further enter the second vacuum chamber 6 via the second aperture 5. The ions are herein subjected to an acceleration and convergence and enter the third aperture 7. The ions, which have entered the MS part 8 via the third aperture 7, are accelerated by an acceleration voltage applied between the ion acceleration electrode 20 and the third aperture 7 as well. The ions then undergo a mass sorting by the quadrupole MS 9. The ions are detected by the detector 10 and provides a mass spectrum after passing through a DC amplifier 11. The first, second and third apertures typically have a skimmer structure, whereby the diffused neutral molecules are prevented from entering the next vacuum chamber. The first vacuum chamber 4 includes no independent vacuum pump and is structured such that this chamber 4 is evacuated by the vacuum pump 1 through the second vacuum chamber 6 from a bypass hole 26 provided downwardly of the second aperture 5. The MS part is evacuated by an independent vacuum pump 2. Numeral 9 designates a quadrupole, and 21 denotes an ion acceleration power supply.

The interface part is heated by the heater power supply 15 and the heater 14 to prevent cooling due to the adiabatic expansion.

Now, it is assumed that the diameters of the first, second and third apertures are 200 μ m, 400 μ m and 500 μ m, respectively; and the diameter of the bypass hole formed downwardly of the second aperture is 5 mm. It is also presumed that the pumping speeds of the vacuum pumps 1, 2 are 16.7 liters/s and 1,000 liters/s. Let P_1 , P_2 , P_3 be the vacuum degrees of the first vacuum chamber, the second vacuum chamber and the MS part. Let C_1 be the conductance of the first aperture 3, and this conductance is defined by the (1) and therefore given as follows:

$$\begin{aligned} C_1 &= 157d^2 \\ &= 157 \times (0.2 \times 10^{-3})^2 \\ &= 6.3 \times 10^{-6} \text{ (m}^3/\text{s)} \end{aligned} \quad (16)$$

Let C_2' be the conductance of the second aperture 5, and let C_2'' be the conductance of the lower bypass hole 26. As $C_2 < C_2''$, the total conductance C_2 from the first vacuum chamber 4 to the second vacuum chamber 6 can be approximated:

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$$\begin{aligned} C_2 &= C_2' + C_2'' \\ &= C_2'' \end{aligned} \quad (17)$$

The conductance in the molecular flow region is given as follows:

$$\begin{aligned} C_2 &= 116 \times 0.834 \times A \\ &= 116 \times 0.834 \times 3.14 \times (2.5 \times 10^{-3})^2 \\ &= 1.9 \times 10^{-3} \text{ (m}^3/\text{s)} \end{aligned} \quad (18)$$

where the coefficient 0.834 is the conductance correction term of the aperture having a thickness.

Assuming that Q_1 is a flow rate of gas flowing into via the first aperture 3 and that Q_2 is a flow rate of gas flowing into the second vacuum chamber 6 from the first vacuum chamber 4, the two flow rates are equal.

$$\begin{aligned} Q_1 &= C_1(P_0 - P_1) \approx C_1 P_0 \\ &= 6.3 \times 10^{-6} \times 10^5 \end{aligned} \quad (19)$$

$$\begin{aligned} &= 0.63 \text{ (Pa} \cdot \text{m}^3/\text{s)} \\ Q_2 &= 1.9 \times 10^{-3} P_1 \end{aligned} \quad (20)$$

As $Q_1 = Q_2$, the pressure P_1 of the first vacuum chamber is given by:

$$\begin{aligned} P_1 &= 0.62 / (1.9 \times 10^{-3}) \\ &= 0.33 \times 10^3 \\ &= 330 \text{ (Pa)} \end{aligned} \quad (21)$$

The pressure P_2 of the second vacuum chamber 6 is given as:

$$\begin{aligned} P_2 &= C_2 P_1 / S_1 \\ &= 1.9 \times 10^{-3} \times 330 / (16.7 \times 10^{-3}) \\ &= 37.5 \text{ (Pa)} \end{aligned} \quad (22)$$

The vacuum obtained in the second vacuum chamber is better than in the first vacuum chamber by approximately one digit. The vacuum P_3 of the MS part is further given as below:

$$\begin{aligned} P_3 &= C_3 P_2 / S_2 \\ &= 116 \times \pi \times (200 \times 10^{-6})^2 \times 38 / (10^3 \times 10^{-3}) \\ &= 5.5 \times 10^{-4} \text{ (Pa)} \end{aligned} \quad (23)$$

This vacuum is enough for the mass analysis.

Parameters of the associated portions under this condition are summarized as follows:

First aperture diameter: 200 μm
 Second aperture diameter: 400 μm
 Third aperture diameter: 500 μm
 Bypass hole diameter: 5 mm
 Pumping speed of pump 1 (e.g., mechanical booster pump): 16.7 liters/s
 Pumping speed of pump 2 (e.g., oil diffusion pump): 1,000 liters/s
 First vacuum chamber pressure: 330 Pa
 Second vacuum chamber pressure: 38 Pa
 MS part vacuum chamber pressure: 5.5×10^{-4} Pa

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When the bypass hole diameter is changed from 5 mm to 2.5 mm, the pressure P_1 of the first vacuum chamber is given as: $330 \times (5/2.5)^2 = 1,320$ Pa. Whereas if changed to 8 mm, the pressure is given as $330 \times (158)^2 = 129$ Pa.

Further, when the number of the bypass hole having a hole diameter of 5 mm is incremented to two, the pressure is given as: $330/2 = 165$ Pa. In this manner, the pressure of the first vacuum chamber can be set simply by changing the bypass hole diameter or the number thereof. In this example, the system is equivalent to the 2-stage differential pumping system shown in the vacuum system diagram of FIG. 8. Namely, the system is equivalent to a 3-stage differential pumping system including an oil rotary pump (pumping speed: 120 liters/m), a mechanical booster pump (pumping speed: 1,000 liters/m) and an oil diffusion pump (pumping speed: 1,000 liters/s). In the interface depicted in FIG. 1, the oil rotary pump, pumping ducts and a vacuum sequence controller are unnecessary, thereby remarkably simplifying the vacuum system.

Assuming that 100 V is applied between the first and second apertures 3, 5, while 10 V is applied between the second and third apertures 5, 7. Further, assuming that the distances between the first, second and third apertures are respectively 5 mm. The pressure of the first vacuum chamber 4 is 330 Pa, while the pressure of the second vacuum chamber 6 is 38 Pa. Hence, the ions are accelerated on the average in the mean free path by an energy of $0.02 \times 20 = 0.4$ (eV) in the first vacuum chamber 4 and by an energy of $0.17 \times 2 = 0.34$ (eV) in the second vacuum chamber 6. Predicted is a spread of an accelerating energy of $0.4 + 0.34 = 0.74$ (eV) at the maximum as a total energy of the two chambers. This value is smaller than 1 eV and falls within such a range as to obtain a sufficient sensitivity and resolving power in either the quadrupole MS or the magnetic sector type MS.

In the first vacuum chamber 4, the ions collide with the neutral molecules (such as nitrogen) 250 times, i.e., $5/0.02 = 250$. With a multiplicity of these collisions, the energy of the collision is converted into an internal energy (equivalent to the heated one) such as vibrations enough to dissociate the added molecules. The highly efficient desolvation is thereby attainable. On the other hand, as illustrated in FIG. 9, in the case of 1-stage differential pumping system, the acceleration by the ion acceleration voltage V_1 of 100 V is effected. When the pressure of the first vacuum chamber 4 is 38 Pa, it follows that a speed spread will be $0.17 \times 100/5 = 3.4$ (eV) at the maximum. The high resolving power and sensitivity can not be obtained any more.

FIG. 10 shows an example where the first vacuum chamber 4 is evacuated only via the second aperture 5. If the diameter of the second aperture is set from several mm to approximately 5 mm, the situation is equivalent to that in the embodiment of FIG. 1.

Another embodiment of the present invention is shown by FIG. 11. Ion sampling is carried out not by the apertures but by a capillary (inside diameter: 0.5–0.2 mm, length: 100 mm–200 mm). The capillary may be made of quartz or metallic material such as stainless steel. In the case of quartz, however, it is required that the ion accelerating electric potential be applicable by effecting silver plating or the like on both ends thereof. Besides, it is possible to help the desolvation by heating this capillary. However, the point that the first vacuum

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chamber is evacuated by the pump 1 via the bypass hole 26 is the same as the embodiment 1.

FIG. 12 shows an insulin (molecular weight: 5734.6) mass spectrum obtained by the conventional system illustrate in FIG. 9. A quantity of introduced sample was 1 μ g. Significant peaks (multiply charged ion) do not appear on the mass spectrum. This measurement involved the use of a double focusing mass spectrometer, wherein the accelerating voltage was 4 kV.

FIG. 13 shows a bovine insulin mass spectrum obtained by the embodiment (FIG. 1) according to the present invention. A quantity of introduced sample was 10 ng. In spite of 1/100 of the above-described sample introduction, there obviously appear insulin's multiply charged ions $(M+6H)^{6+}$, $(M+5H)^{5+}$, and $(M+4H)^{4+}$. It can be considered that the desolvation was imperfect in the foregoing system, and the multiply charged ions irregularly appear as noises in a wide mass region or captured by the electric field of the double focusing mass spectrometer. In accordance with the embodiment of the present invention, the desolvation of the multiply charged ions was sufficiently performed, and the mass peak is obviously given onto the mass spectrum. Further, the noises on the mass spectrum due to the cluster ions are reduced.

As discussed above, the multiply charged ions and pseudomolecular ions are subjected the sufficient desolvation, and the measurement can be performed with a high sensitivity.

The electro-spray ionization (ESI) has been exemplified as the atmospheric pressure ionization. The same effects are, however, obtainable by atmospheric pressure chemical ionization (APCI), pneumatically assisted ESI and the like. Further, the present invention is applied not only to the LC/MS but to methods of ionization under the atmospheric pressure as in the case of supercritical fluid chromatography (SFC)/MS and CZE (Capillary Zone Electrophoresis)/MS.

According to the present invention, the differential pumping system is simplified, and the inexpensive device can be provided.

What is claimed is:

1. A mass spectrometry comprising the steps of: generating ions under an atmospheric pressure, evacuating low and intermediate vacuum chambers by a common vacuum system so that a vacuum of the former is lower than that of the latter, leading the generated ions to a high vacuum chamber having a vacuum higher than in said low and intermediate vacuum chambers via said low and intermediate vacuum chambers, and effecting a mass analysis therein.

2. The mass spectrometry as set forth in claim 1, wherein said low vacuum chamber is evacuated by said vacuum system through said intermediate vacuum chamber.

3. The mass spectrometry as set forth in claim 2, further comprising a step of accelerating the ions respectively by a first accelerating voltage in said low vacuum chamber and by an accelerating voltage lower than the first accelerating voltage in said intermediate vacuum chamber.

4. The mass spectrometry as set forth in claim 1, further comprising a step of accelerating the ions respectively by a first accelerating voltage in said low vacuum chamber and by an accelerating voltage lower than the first accelerating voltage in said intermediate vacuum chamber.

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5. A mass spectrometer comprising: a means for generating ion under an atmospheric pressure; a means for effecting a mass analysis of the ions under a high vacuum; and an interface disposed between said two means to connect said two means, characterized in that: said interface includes a low vacuum chamber, an intermediate vacuum chamber disposed between said low vacuum chamber and said mass analysis means and a means for evacuating said chambers so that a vacuum of the former is lower than in the latter; said evacuating means is common to said low and intermediate vacuum chambers; and said low and intermediate vacuum chambers include openings through which the ions generated by said ion generating means pass towards said mass analysis means.

6. The mass spectrometer as set forth in claim 5, wherein at least one of said openings has a skimmer configuration.

7. The mass spectrometer as set forth in claim 6, further comprising a bypass pumping hole through which said low vacuum chamber communicates with said intermediate vacuum chamber so that said low vacuum chamber is evacuated by said evacuating means through said intermediate vacuum chamber.

8. The mass spectrometer as set forth in claim 7, further comprising a means for accelerating the ions respectively by a first accelerating voltage in said low vacuum chamber and by an accelerating voltage lower than the first accelerating voltage in said intermediate vacuum chamber.

9. The mass spectrometer as set forth in claim 5, further comprising a bypass pumping hole through which said low vacuum chamber communicates with said intermediate vacuum chamber so that said low vacuum chamber is evacuated by said evacuating means through said intermediate vacuum chamber.

10. The mass spectrometer as set forth in claim 9, further comprising a means for accelerating the ions respectively by a first accelerating voltage in said low vacuum chamber and by an accelerating voltage lower than the first accelerating voltage in said intermediate vacuum chamber.

11. The mass spectrometer as set forth in claim 5, further comprising a means for accelerating the ions respectively by a first accelerating voltage in said low vacuum chamber and by an accelerating voltage lower than the first accelerating voltage in said intermediate vacuum chamber.

12. An interface comprising: first and second vacuum chambers; and a means for evacuating said vacuum chambers so that a vacuum of the former is lower than in the latter, characterized in that said first and second vacuum chambers include openings disposed so that the ions are allowed to pass through said first and second vacuum chambers, and said evacuating means is common to said first and second vacuum chambers.

13. The interface as set forth in claim 12, wherein said openings are formed in a skimmer configuration.

14. The interface as set forth in claim 13, further comprising a means for accelerating the ions by a first accelerating voltage in said first vacuum chamber and by a second accelerating voltage lower than said first accelerating voltage in said second vacuum chamber.

15. The interface as set forth in claim 12, further comprising a means for accelerating the ions by a first accelerating voltage in said first vacuum chamber and by a second accelerating voltage lower than said first accelerating voltage in said second vacuum chamber.

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16. A mass spectrometer comprising:
a means for generating ions under an atmospheric
pressure; a first vacuum chamber disposed to pass
the ions therethrough; a second vacuum chamber
disposed to pass therethrough the ions which have
been passed through said first vacuum chamber; a
means for mass-analyzing the ions which have
passed through said second vacuum chamber; a
means for evacuating said first and second vacuum

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chambers so that a vacuum of the former is lower
than in the latter;
a means for accelerating the ions by a first accelerat-
ing voltage in said first vacuum chamber and by an
accelerating voltage lower than said first accelerat-
ing voltage in said second vacuum chamber; and
a bypass pumping hole through which said first vac-
uum chamber communicates with said second vac-
uum chamber so that said first vacuum chamber is
evacuated by said evacuating means through said
second vacuum chamber.

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CERTIFICATE OF SERVICE

I, Karen E. Keller, Esquire hereby certify that on December 15, 2005, I caused to be electronically filed a true and correct copy of the foregoing document with the Clerk of the Court using CM/ECF, which will send notification that such filing is available for viewing and downloading to the following counsel of record:


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